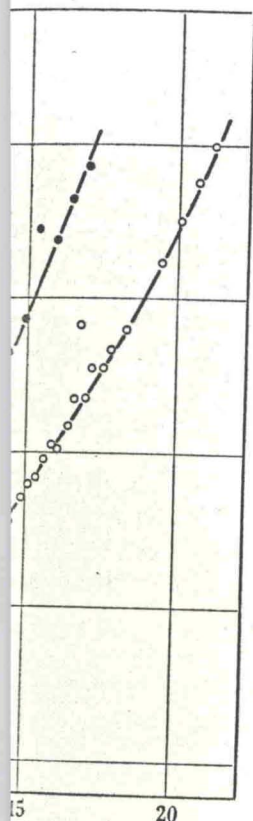


is also possible to treat the results at various pressures and temperatures so that they are not so good. Such discrepancy between our values of  $T/\phi < 0.10$  and our Debye  $\theta$  values while Keesom & Keesom's values



increase with falling temperature. We may notice that Webb, Wilkinson & Wilks (1952) have also observed this effect.

In view of this relationship it is convenient to present certain thermodynamic properties of the solid in terms of  $T/\phi$  and table 4 gives  $C_v$ ,  $S$  and  $(U - U_0)/T$  as

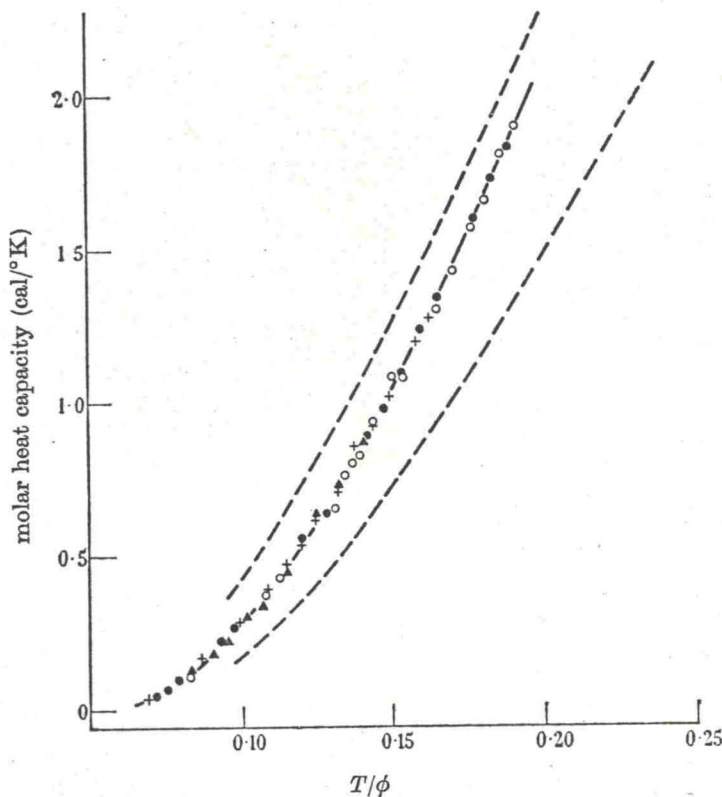


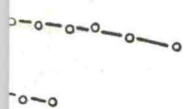
FIGURE 8. The molar-heat capacities of figure 6 plotted against the reduced temperature  $T/\phi$ .  $\circ$ , 10.6 ml.;  $\bullet$ , 11.7 ml.; +, 13.0 ml.;  $\blacktriangle$ , 14.4 ml.

TABLE 3. SMOOTHED VALUES OF THE PARAMETER  $\phi$  AS A FUNCTION OF VOLUME

V (ml.)	$\phi$	V (ml.)	$\phi$
10.5	113	15	48.5
11	101.5	16	41.4
12	83.2	17	35.7
13	68.7	18	31.1
14	57.4	19	27.2
		20	24.2

functions of  $T/\phi$ . (The units of  $\phi$  are arbitrary; they have, however, been chosen so as to be similar in magnitude to the Debye  $\theta$ 's.) These values depend on an extrapolation of the specific heats to  $0^\circ\text{K}$ . This extrapolation entails some uncertainty owing to the above-mentioned discrepancy between our results and those of Keesom & Keesom. As, however, the maximum error will only be of the order of 0.01 entropy unit, it is quite immaterial for our purposes. For a fuller discussion of this point we refer to a paper by Webb & Wilks (1953).

of solid helium.  
ml.;  $\blacktriangle$ , 14.4 ml.



20

t four molar volumes.